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Harmonic Maps into Principal Bundles and Generalized Magnetic Mappings

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Outline

- Motivation
- Harmonic maps into Kaluza-Klein principal bundles
 - Kaluza-Klein metrics on principal bundles
 - Kaluza-Klein harmonic maps and generalized magnetic maps
- Examples of Kaluza-Klein harmonic maps and generalized magnetic maps

Source

The results presented in this talk are based on the following paper:

H. Benziadi, A. López Almorox, and C. Tejero Prieto. “Harmonic maps into principal bundles and generalized magnetic mappings”. In: *Journal of Geometry and Physics* (2026)

Motivation

Riemannian magnetic curves

Let $\gamma : I \rightarrow (M, g)$ be a smooth curve, and let $A \in \Omega^1(M)$ be a gauge potential. Its curvature is a closed 2-form

$$F = dA,$$

called the (**Faraday**) *magnetic field*.

The associated Lorentz endomorphism $\mathcal{F} \in T_1^1(M)$ is defined by

$$g(\mathcal{F}(D_1), D_2) = F(D_1, D_2), \quad \forall D_1, D_2 \in \mathfrak{X}(M).$$

A smooth curve γ is called a *magnetic curve* if it satisfies the Lorentz equation

$$\nabla_{\dot{\gamma}} \dot{\gamma} = q \mathcal{F}(\dot{\gamma}),$$

where $q \in \mathbb{R}$ is a constant representing the charge, and $\mathcal{F}(\dot{\gamma})$ is the Lorentz force.

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Kaluza-Klein formalism for Riemannian magnetic curves

Let $P = M \times U(1)$ be a $U(1)$ -principal bundle over M with a principal connection $\omega \in \Omega^1(P, \mathfrak{u}(1))$, where $\dim M = m$ and $\dim P = m + 1$.

We consider the commutative diagram

$$\begin{array}{ccc} I \subset \mathbb{R} & \xrightarrow{\tilde{\gamma}} & (P = M \times U(1), \hat{g}) \\ & \searrow \gamma & \downarrow \pi, G=U(1) \\ & & (M, g) \end{array}$$

The metric $\hat{g}$ is called the *Kaluza-Klein metric* and is defined by

$$\hat{g} = \pi^* g + \kappa \langle \omega, \omega \rangle,$$

where $\kappa > 0$ is a constant and $\langle \cdot, \cdot \rangle$ is an Ad-invariant inner product on $\mathfrak{u}(1)$.

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Using the global section $s \in \Gamma(M, P)$, the Kaluza–Klein metric is represented by

$$\hat{g}_{AB} = \begin{pmatrix} g_{\mu\nu} + \kappa^2 A_\mu A_\nu & \kappa A_\nu \\ \kappa A_\mu & 1 \end{pmatrix}, \quad s^* \omega = A \in \Omega^1(M, \mathfrak{g}).$$

Let $\tilde{\gamma} : I \longrightarrow (P, \hat{g})$ be a geodesic curve on P , that is:

$$\frac{d^2 x^K}{dt^2} + \Gamma_{BC}^K \frac{dx^B}{dt} \frac{dx^C}{dt} = 0, \quad K, B, C = 1, \dots, m+1.$$

Its projection γ on the base manifold M satisfies:

$$\frac{d^2 x^\mu}{dt^2} + \Gamma_{\nu\sigma}^\mu \frac{dx^\nu}{dt} \frac{dx^\sigma}{dt} = \kappa \left(\frac{dx^{m+1}}{dt} + \kappa A_\nu \frac{dx^\nu}{dt} \right) \mathcal{F}_\nu^\mu \frac{dx^\nu}{dt}, \quad \mu, \nu, \sigma = 1, \dots, m \quad (1)$$

where $\mathcal{F}_\nu^\mu = g^{\mu\sigma} F_{\sigma\nu}$ and $\mathcal{F}_\nu^\mu \frac{dx^\nu}{dt}$ represents the Lorentz force.

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Therefore,

$$\frac{dx^{m+1}}{dt} + \kappa A_\nu \frac{dx^\nu}{dt} = \text{constant}.$$

Hence, we obtain

$$\frac{d^2 x^\mu}{dt^2} + \Gamma_{\nu\sigma}^\mu \frac{dx^\nu}{dt} \frac{dx^\sigma}{dt} = \bar{\kappa} \mathcal{F}^\mu{}_\nu \frac{dx^\nu}{dt}.$$

where

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Generalization to Arbitrary Gauge Groups G : Kerner's Generalization

$$\begin{array}{ccc} I \subset \mathbb{R} & \xrightarrow{\tilde{\gamma}} & P = M \times G \\ & \searrow \gamma & \downarrow \pi \\ & & (M, g) \end{array}$$

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Harmonic maps into Kaluza-Klein principal bundles

To extend the Kaluza-Klein formalism beyond curves, we replace the interval $(I \in \mathbb{R})$ by a Riemannian manifold (N, h) with $\dim N > 1$. Since geodesics are harmonic maps when the domain is one-dimensional, harmonic maps provide the natural higher-dimensional generalization of geodesics.

We consider the following commutative diagram:

$$\begin{array}{ccc} (N, h) & \xrightarrow{\tilde{\Phi}} & (P, \hat{g}) \\ & \searrow \Phi & \downarrow \pi, G \\ & & (M, g) \end{array}$$

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$$\dim M = m, \quad \dim G = d, \quad \dim P = m + d.$$

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Harmonic maps

- Let (N, h) and (M, g) be two Riemannian manifolds, with N compact and oriented.
- Let $\Phi : N \longrightarrow M$ be a smooth map.
- We define the Dirichlet energy functional

$$E(\Phi) = \frac{1}{2} \int_N \|d\Phi\|^2 \text{Vol}_h .$$

- A smooth map $\Phi : N \longrightarrow M$ is said to be harmonic if it is a critical point of $E(\Phi)$, i.e., it satisfies the Euler-Lagrange equation:

$$\tau(\Phi) = \text{Tr}_h(\nabla d\Phi) = 0,$$

where $\Pi_\Phi := \nabla d\Phi \in S^2(N, \Phi^*TM)$ and $\tau(\Phi) \in \Gamma(N, \Phi^*TM)$.

- In terms of local coordinates $\{y^s\}$ on N and $\{x^\mu\}$ on M , we have

$$g^{st} \left[\frac{\partial^2 \Phi^\lambda}{\partial y^s \partial y^t} - \Gamma_{st}^l \frac{\partial \Phi^\lambda}{\partial y^l} + \frac{\partial \Phi^\mu}{\partial y^t} \frac{\partial \Phi^\nu}{\partial y^s} (\bar{\Gamma}_{\mu\nu}^\lambda \circ \Phi) \right] = 0, \quad \lambda = 1, \dots, m.$$

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Special case: Geodesics ($\dim N = 1$)

For $N = I \subset \mathbb{R}$, a harmonic map $\gamma : I \rightarrow (M, g)$ is precisely a geodesic satisfying

$$\frac{d^2 x^\mu}{dt^2} + \Gamma_{\nu\sigma}^\mu \frac{dx^\nu}{dt} \frac{dx^\sigma}{dt} = 0.$$

Kaluza-Klein metrics on principal G -bundles

Definition 1

A Kaluza-Klein metric $\widehat{g}$ on a principal G -bundle $\pi : P \longrightarrow M$ is a Riemannian metric on P that is G -invariant.

Theorem 1

The space $\text{Met}_{KK}(P)$ of Kaluza-Klein metrics on P is equipped with an $\text{Aut}(P)$ -equivariant bijection

$$\begin{aligned}\Psi : \text{Met}_{KK}(P) &\xrightarrow{\sim} \mathcal{A}(P) \times \text{Met}(\text{ad}P) \times \text{Met}(M) \\ \widehat{g} &\longmapsto (\omega, \quad g_{\text{ad}}, \quad g)\end{aligned}$$

$$TP = \mathcal{V}P \oplus (\mathcal{V}P)^\perp, \quad \pi : (P, \widehat{g}) \rightarrow (M, g) \text{ is a Riemannian submersion.}$$

The inverse bijection is given by

$$\widehat{g} := \Psi^{-1}(\omega, g_{\text{ad}}, g) = (\omega^* g_{\text{ad}})^\mathcal{V} + \pi^* g.$$

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Kaluza-Klein harmonic maps and generalized magnetic maps

Let $\pi : (P, \widehat{g}) \longrightarrow (M, g)$ be a Kaluza-Klein principal G -bundle

$$\begin{array}{ccc} (N, h) & \xrightarrow{\widetilde{\Phi}} & (P, \widehat{g}) \\ & \searrow \Phi & \downarrow \pi, G \\ & & (M, g) \end{array}$$

We establish the relationship between the tension fields of $\widetilde{\Phi}$ and Φ .

Proposition 1

The second fundamental form $\Pi_\Phi \in S^2(N, \Phi^*TM)$ of the composition $\Phi = \pi \circ \tilde{\Phi} : (N, h) \longrightarrow (M, g)$ satisfies

$$\Pi_\Phi = \pi_* \circ \Pi_{\tilde{\Phi}} + \tilde{\Phi}^*(\Pi_\pi)$$

and its h -Riemannian trace satisfies the identity

$$\tau(\Phi) = \pi_*(\tau(\tilde{\Phi})) + \text{Tr}_h [\tilde{\Phi}^*(\Pi_\pi)].$$

Corollary 1

If $\tilde{\Phi}$ is harmonic, then

$$\tau(\Phi) = \text{Tr}_h [\tilde{\Phi}^*(\Pi_\pi)].$$

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Proposition 2

Let $\pi : (P, \widehat{g}) \longrightarrow (M, g)$ be a Kaluza-Klein principal G -bundle. For every π -projectable vector field $D \in \mathfrak{X}(P)$, we have

$$\Pi_\pi(D, D) := (\nabla d\pi)(D, D) = [-\pi_* \circ (2\mathbf{A} + \mathbf{T})](D, D),$$

where ∇ is the connection on $T^*P \otimes \pi^*TM$ induced by the Levi-Civita connections of $(P, \widehat{g})$ and (M, g) , $\mathbf{A}$ and $\mathbf{T}$ are the O'Neill's tensors defined by

$$\mathbf{A}(D_1, D_2) = [\nabla_{D_1^{\mathcal{H}}}^P D_2^{\mathcal{V}}]^{\mathcal{H}} + [\nabla_{D_1^{\mathcal{H}}}^P D_2^{\mathcal{H}}]^{\mathcal{V}},$$

$$\mathbf{T}(D_1, D_2) = [\nabla_{D_1^{\mathcal{V}}}^P D_2^{\mathcal{H}}]^{\mathcal{V}} + [\nabla_{D_1^{\mathcal{V}}}^P D_2^{\mathcal{V}}]^{\mathcal{H}}.$$

Theorem 2

Let $\pi : (P, \widehat{g}) \rightarrow (M, g)$ be a Kaluza–Klein principal G -bundle, $\widetilde{\Phi} : (N, h) \rightarrow (P, \widehat{g})$ a smooth map, and let $\Phi = \pi \circ \widetilde{\Phi} : (N, h) \rightarrow (M, g)$. Then the tension field of Φ is given by

$$\tau(\Phi) = \pi_*(\tau(\widetilde{\Phi})) - \pi_* \left(\text{Tr}_h [\widetilde{\Phi}^*(2\mathbf{A} + \mathbf{T})] \right).$$

In particular, $\widetilde{\Phi}$ is (horizontally) harmonic if and only if

$$\tau(\Phi) = -\pi_* \left(\text{Tr}_h [\widetilde{\Phi}^*(2\mathbf{A})] \right) - \pi_* \left(\text{Tr}_h [\widetilde{\Phi}^*\mathbf{T}] \right).$$

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$$\tau(\Phi) = -\pi_* \left(\text{Tr}_h [\widetilde{\Phi}^*(2\mathbf{A})] \right) - \pi_* \left(\text{Tr}_h [\widetilde{\Phi}^*\mathbf{T}] \right).$$

Let $\pi : (P, \widehat{g}) \longrightarrow (M, g)$ be a Kaluza-Klein principal G -bundle.

Proposition 3

The vertical component of the tension field of a smooth map $\widetilde{\Phi} : (N, h) \longrightarrow (P, \widehat{g})$ is given by

$$[\tau(\widetilde{\Phi})]^\nu = -\delta^{(\nabla^N, \widetilde{\Phi}^* \nabla^{P, \nu})}(\widetilde{\Phi}^* \omega) + \omega(\text{Tr}_h(\widetilde{\Phi}^* \mathbf{T})).$$

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Lorentz Endomorphism

Let $\pi : (P, \widehat{g}) \longrightarrow (M, g)$ be a Kaluza-Klein principal bundle with a principal connection $\omega \in \Omega^1(P, \mathcal{V}P)$. Its curvature $\Omega^\omega \in \Omega^2(P, \mathcal{V}P)$ induces the **Lorentz endomorphism**

$\mathcal{F}^\omega \in T_1^1(P) \otimes_{C^\infty(P)} \mathfrak{X}^\mathcal{V}(P)$, defined by

$$\langle V, \Omega^\omega \rangle_{\mathcal{V}P}(D_1, D_2) = \widehat{g}(\langle V, \mathcal{F}^\omega \rangle(D_1), D_2),$$

for every vertical vector field V and all $D_1, D_2 \in \mathfrak{X}(P)$

Definition 2

The curvature-modified metric $\widehat{g}_{\mathcal{F}^\omega}$ is defined by

$$\widehat{g}_{\mathcal{F}^\omega}(D_1, D_2) = \frac{1}{2} \{ \langle D_1^\mathcal{V}, \mathcal{F}^\omega \rangle_{\mathcal{V}P}(D_2^\mathcal{H}) + \langle D_2^\mathcal{V}, \mathcal{F}^\omega \rangle_{\mathcal{V}P}(D_1^\mathcal{H}) \}.$$

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$$\tau(\Phi) = -\pi_* \left(\text{Tr}_h [\widetilde{\Phi}^* \widehat{g}_{\mathcal{F}^\omega}] \right) - \pi_* \left(\text{Tr}_h [\widetilde{\Phi}^* \mathbf{T}] \right).$$

Definition 3

The equation above is called the **intrinsic generalized Lorentz equation** satisfied by the **generalized magnetic map** Φ .

Definition 4

A generalized magnetic map Φ is called **uncharged** if

$$\pi_* \left(\text{Tr}_h [\widetilde{\Phi}^* \widehat{g}_{\mathcal{F}^\omega}] \right) = 0.$$

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Theorem 5

Let $\tilde{\Phi} : (N, h) \longrightarrow (P, \hat{g})$ be a smooth map and $\Phi = \pi \circ \tilde{\Phi}$. Then $\tilde{\Phi}$ is harmonic if and only if

$$\tau(\Phi) = -\pi_*[\mathrm{Tr}_h(\tilde{\Phi}^*\hat{g}_{\mathcal{F}\omega})] - \pi_*[\mathrm{Tr}_h(\tilde{\Phi}^*\mathbf{T})] \quad (2)$$

and

$$\delta^{(\nabla^N, \tilde{\Phi}^*\nabla^{P, \mathcal{V}})}(\tilde{\Phi}^*\omega) = \omega[\mathrm{Tr}_h(\tilde{\Phi}^*\mathbf{T})]. \quad (3)$$

These equations are called **generalized Wong's equations**

Definition 5

A **Kaluza-Klein harmonic map** is a harmonic map whose target is a Kaluza-Klein principal bundle.

Theorem 5

Let $\tilde{\Phi} : (N, h) \rightarrow (P, \hat{g})$ be a smooth map and $\Phi = \pi \circ \tilde{\Phi}$. Then $\tilde{\Phi}$ is harmonic if and only if

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A **Kaluza-Klein harmonic map** is a harmonic map whose target is a Kaluza-Klein principal bundle.

Let $\pi : (P, \widehat{g}) \longrightarrow (M, g)$ be a Kaluza-Klein principal G -bundle with a principal connection $\omega \in \Omega^1(P, \mathcal{V}P)$. Its curvature $\Omega^\omega \in \Omega^2(P, \mathcal{V}P)$ is horizontal and G -invariant.

Lorentz strength endomorphism

Let $B = \{\xi_1, \dots, \xi_d\}$ be a basis of the Lie algebra $\mathfrak{g}$.

Since the curvature form Ω^ω is G -invariant, it descends to

$$\Omega_\omega = \sum_{a=1}^d \Omega_\omega^a \otimes \xi_a^* \in \Omega^2(M) \otimes_{C^\infty(M)} \mathfrak{X}^\mathcal{V}(P)^G,$$

which induces the Lorentz strength endomorphism

$$\mathcal{F}_\omega = \sum_{a=1}^d \mathcal{F}_\omega^a \otimes \widetilde{\xi}_a \in T_1^1(M) \otimes_{C^\infty(M)} \mathfrak{X}^\mathcal{V}(P)^G.$$

Theorem 6

The map $\tilde{\Phi} : (N, h) \longrightarrow (P, \hat{g})$ is (horizontally) harmonic if and only if $\Phi = \pi \circ \tilde{\Phi}$ satisfies

$$\tau(\Phi) = - \sum_{a=1}^k [\Phi^*(\mathcal{F}_\omega^a)] \left(p_h^{-1}(\tilde{\Phi}^*[p_{\hat{g}}(\tilde{\xi}_a)]) \right) - \pi_* \left(\text{Tr}_h [\tilde{\Phi}^* \mathbf{T}] \right)$$

Corollary 2

If, for each a , the vector field $\tilde{\xi}_a$ is a Killing vector field, then

$$K_a := p_h^{-1}(\tilde{\Phi}^*[p_{\hat{g}}(\tilde{\xi}_a)]) \in \mathfrak{X}(N)$$

represents the **conserved current** along $\tilde{\Phi}$.

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We now introduce the natural generalization of the classical example of geodesic curves.

Example 1: Generalized magnetic curves

Let $\tilde{\gamma} : (I, h) \longrightarrow (P, \hat{g})$ be a geodesic curve in P , that is, $\nabla_{\dot{\tilde{\gamma}}}^P \dot{\tilde{\gamma}} = 0$, and let $\gamma = \pi \circ \tilde{\gamma}$ the induced curve in M , then $\tilde{\gamma}$ satisfies

$$\begin{aligned} \nabla_{\dot{\gamma}}^M \dot{\gamma} &= -\pi_*[\mathrm{Tr}_h(\tilde{\gamma}^* \hat{g}_{\mathcal{F}\omega})] - \pi_*[\mathrm{Tr}_h(\tilde{\gamma}^* \mathbf{T})], \\ \delta^{(\nabla^I, \tilde{\gamma}^* \nabla^{P, \mathcal{V}})}(\tilde{\gamma}^* \omega) &= \omega[\mathrm{Tr}_h(\tilde{\gamma}^* \mathbf{T})]. \end{aligned}$$

The first equation is equivalent to

$$\nabla_{\dot{\gamma}}^M \dot{\gamma} = - \sum_{a=1}^k \tilde{\gamma}^\# \widehat{g}(\tilde{\gamma}^\# \tilde{\xi}_a, \dot{\gamma}) \mathcal{F}_\omega^a(\dot{\gamma}) - \pi_*[\mathbf{T}(\dot{\gamma}, \dot{\gamma})] \quad (4)$$

Suppose that, for each a , the vector field $\tilde{\xi}_a$ is a Killing vector field with respect to $\widehat{g}$ on P . It follows that

$$\kappa_a = -\tilde{\gamma}^\# \widehat{g}(\tilde{\gamma}^\# \tilde{\xi}_a, \dot{\gamma})$$

is constant. Therefore,

$$\begin{aligned} \nabla_{\dot{\gamma}}^M \dot{\gamma} &= \sum_{a=1}^d \kappa_a \cdot \mathcal{F}_\omega^a(\dot{\gamma}) - \pi_*[\mathbf{T}(\dot{\gamma}, \dot{\gamma})], \\ \delta^{(\nabla^I, \tilde{\gamma}^* \nabla^{P, \mathcal{V}})}(\tilde{\gamma}^* \omega) &= \omega[\mathbf{T}(\dot{\gamma}, \dot{\gamma})]. \end{aligned}$$

The first equation is equivalent to

$$\nabla_{\dot{\gamma}}^M \dot{\gamma} = - \sum_{a=1}^k \tilde{\gamma}^{\#} \widehat{g}(\tilde{\gamma}^{\#} \tilde{\xi}_a, \dot{\gamma}) \mathcal{F}_{\omega}^a(\dot{\gamma}) - \pi_*[\mathbf{T}(\dot{\gamma}, \dot{\gamma})] \quad (4)$$

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Example 2

If $N = P$, $g = \widehat{g}$, then $\text{Id}_P : (P, \widehat{g}) \rightarrow (P, \widehat{g})$ is harmonic.

$$\begin{array}{ccc} (P, \widehat{g}) & \xrightarrow{\text{Id}_P} & (P, \widehat{g}) \\ & \searrow \pi & \downarrow \pi \\ & & (M, g) \end{array}$$

The Wong's equations satisfied by Id_P are

$$\tau(\pi) = -d \cdot \pi_*(\mathbf{H}^\mathcal{V}), \quad \delta^{(\nabla^N, \nabla^\omega)} \omega = \mathbf{V},$$

with $\mathbf{H}^\mathcal{V} := \frac{1}{d} \text{Tr}_{\widehat{g}} \mathbf{T}^{\mathcal{V}, \mathcal{V}}$ the vertical mean curvature vector field of the fibers of π and $\mathbf{V} = \left[\sum_{a=1}^d a d_{\xi_a}^* (\xi_a) \right]^*$, where $\{\xi_a^*\}_{a=1}^d$ is a frame $\widehat{g}$ -ortonormal of $\mathcal{V}P$.

Example 3

Let

$$\tilde{\Phi}_\alpha : S^1 \times S^1 \longrightarrow S^3 \subset \mathbb{R}^4$$

be the map defined by

$$\tilde{\Phi}_\alpha(z_1, z_2) = \cos \alpha z_1 + \sin \alpha z_2 \cdot j.$$

The following diagram is commutative:

$$\begin{array}{ccc} S^1 \times S^1 & \xrightarrow{\tilde{\Phi}_\alpha} & S^3 \subset \mathbb{R}^4 \\ & \searrow \Phi_\alpha & \downarrow \pi \\ & & S^2 \subset \mathbb{R}^3 \end{array}$$

where $\Phi_\alpha : S^1 \times S^1 \longrightarrow S^2 \subset \mathbb{R}^3$ is the composition $\Phi_\alpha = \pi \circ \tilde{\Phi}_\alpha$.

The vertical and horizontal subbundles of the Hopf $U(1)$ -bundle (S^3, R) are given by $V(S^3) = \langle i^* \rangle$, $H(S^3) = \langle j^*, k^* \rangle$.

The principal connection is given by $\omega = \omega^1 \otimes i^*$.

The curvature of the connection of ω is

$$\Omega^\omega = d\omega^1 \otimes i^* = -2\omega^2 \wedge \omega^3 \otimes i^*.$$

The Lorentz strength endomorphism of $(S^3, \widehat{g})$ is

$$\mathcal{F}^\omega = -2(\omega^2 \otimes D_3 - \omega^3 \otimes D_2) \otimes i^*.$$

The curvature-modified metric of $\hat{g}$ is

$$\begin{aligned}\hat{g}_{\mathcal{F}\omega} &= -[\omega^1 \otimes \omega^2 + \omega^2 \otimes \omega^1] \otimes D_3 + [\omega^1 \otimes \omega^3 + \omega^3 \otimes \omega^1] \otimes D_2 \\ &= \omega^1 \cdot \omega^3 \otimes D_2 - \omega^1 \cdot \omega^2 \otimes D_3\end{aligned}$$

The Wong's equations satisfied by $\tilde{\Phi}$ are

$$\tau(\Phi_\alpha) = -\pi_* \left(\text{Tr}_g \left[\tilde{\Phi}_\alpha^* (-\omega^1 \cdot \omega^2 \otimes D_3 + \omega^1 \cdot \omega^3 \otimes D_2) \right] \right)$$

$$\delta^{(\nabla^{S^1 \times S^1}, \tilde{\Phi}_\alpha^* \nabla^\omega)} (\tilde{\Phi}_\alpha^* (\omega^1 \otimes i^*)) = 0$$

If

$$\alpha \neq \frac{\pi}{2}\mathbb{Z},$$

then $\tilde{\Phi}_\alpha$ is a spherical harmonic immersion, called the α -twisted Clifford torus; that is,

$$\tau(\tilde{\Phi}_\alpha) = 0$$

(in particular, for $\alpha = \frac{\pi}{4}$). Hence, Φ_α is a generalized magnetic map.

A straightforward computation shows that Φ_α is **uncharged** (i.e., $\tau(\Phi_\alpha) = 0$) if and only if

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Example 4

Let

$$\tilde{\Phi}_\alpha : S^3 \times S^3 \longrightarrow S^7 \subset \mathbb{R}^8$$

be a map given by

$$\tilde{\Phi}_\alpha(q_1, q_2) = \cos\alpha q_1 + \sin\alpha q_2 \cdot l.$$

We have the following commutative diagram:

$$\begin{array}{ccc} S^3 \times S^3 & \xrightarrow{\tilde{\Phi}_\alpha} & S^7 \subset \mathbb{R}^8 \\ & \searrow \Phi_\alpha & \downarrow \pi \\ & & S^4 \subset \mathbb{R}^5 \end{array}$$

where $\Phi_\alpha : S^3 \times S^3 \longrightarrow S^4 \subset \mathbb{R}^5$ is the composition $\Phi_\alpha = \pi \circ \tilde{\Phi}_\alpha$.

The vertical and horizontal subbundles of the Hopf $SU(2)$ -bundle (S^7, R) are

$$V(S^7) = \langle D_1 = \mathbf{i}^*, D_2 = \mathbf{j}^*, D_3 = \mathbf{k}^* \rangle,$$

$$H(S^7) = \langle D_4 = \mathbf{l}^*, D_5 = \mathbf{m}^*, D_6 = \mathbf{n}^*, D_7 = \mathbf{o}^* \rangle.$$

The principal connection is

$$\omega = \omega^1 \otimes \mathbf{i}^* + \omega^2 \otimes \mathbf{j}^* + \omega^3 \otimes \mathbf{k}^*.$$

The Curvature of the connection ω is

$$\begin{aligned} \Omega^\omega = & 2[(-\omega^4 \wedge \omega^5 + \omega^6 \wedge \omega^7) \otimes \mathbf{i}^* - (\omega^4 \wedge \omega^6 + \omega^5 \wedge \omega^7) \otimes \mathbf{j}^* + \\ & + (-\omega^4 \wedge \omega^7 + \omega^5 \wedge \omega^6) \otimes \mathbf{k}^*]. \end{aligned}$$

The Lorentz strength endomorphism of $(S^7, \widehat{g})$ is

$$\begin{aligned}\mathcal{F}^\omega = & 2[(-\omega^4 \otimes D_5 + \omega^5 \otimes D_4 + \omega^6 \otimes D_7 - \omega^7 \otimes D_6) \otimes \mathbf{i}^* + \\ & - (\omega^4 \otimes D_6 - \omega^6 \otimes D_4 + \omega^5 \otimes D_7 - \omega^7 \otimes D_5) \otimes \mathbf{j}^* + \\ & + (-\omega^4 \otimes D_7 + \omega^7 \otimes D_4 + \omega^5 \otimes D_6 - \omega^6 \otimes D_5) \otimes \mathbf{k}^*].\end{aligned}$$

The curvature-modified metric of $\widehat{g}$ is

$$\begin{aligned}\widehat{g}_{\mathcal{F}^\omega} = & -\omega^1 \cdot \omega^4 \otimes D_5 + \omega^1 \cdot \omega^5 \otimes D_4 + \omega^1 \cdot \omega^6 \otimes D_7 + \omega^1 \cdot \omega^7 \otimes D_6 - \\ & - \omega^2 \cdot \omega^4 \otimes D_6 - \omega^2 \cdot \omega^5 \otimes D_7 + \omega^2 \cdot \omega^6 \otimes D_4 + \omega^2 \cdot \omega^7 \otimes D_5 - \\ & - \omega^3 \cdot \omega^4 \otimes D_7 + \omega^3 \cdot \omega^5 \otimes D_6 - \omega^3 \cdot \omega^6 \otimes D_5 + \omega^3 \cdot \omega^7 \otimes D_4.\end{aligned}$$

The Wong's equations satisfied by $\tilde{\Psi}_\alpha$ are

$$\begin{aligned}\tau(\Psi_\alpha) = & -\pi_* \left(\text{Tr}_g [\tilde{\Psi}_\alpha ((\omega^1 \cdot \omega^5 + \omega^2 \cdot \omega^6 + \omega^3 \cdot \omega^7) \otimes D_4 - \right. \\ & - (\omega^1 \cdot \omega^4 - \omega^2 \cdot \omega^7 + \omega^3 \cdot \omega^6) \otimes D_5 + \\ & + (\omega^1 \cdot \omega^7 - \omega^2 \cdot \omega^4 + \omega^3 \cdot \omega^5) \otimes D_6 + \\ & \left. - (\omega^1 \cdot \omega^6 - \omega^2 \cdot \omega^5 - \omega^3 \cdot \omega^4) \otimes D_7) \right])\end{aligned}$$

and

$$\delta(\nabla^{S^3 \times S^3}, \tilde{\Psi}_\alpha^* \nabla^\omega)(\tilde{\Psi}_\alpha^*(\omega^1 \otimes \mathbf{i}^* + \omega^2 \otimes \mathbf{j}^* + \omega^3 \otimes \mathbf{k}^*)) = 0.$$

If

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Thank you for your attention!

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